

Locality and Nonlocality in Quantum Pure-state Identification Problems

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Suppose we want to identify an input state with one of two unknown reference states, where the input state is guaranteed to be equal to one of the reference states. We assume that no classical knowledge of the reference states is given, but a certain number of copies of them are available instead. Two reference states are independently and randomly chosen from the state space in a unitary invariant way. This is called the quantum state identification problem, and the task is to optimize the mean identification success probability. In this paper, we consider the case where each reference state is pure and bipartite, and generally entangled. The question is whether the maximum mean identification success probability can be attained by means of a local operations and classical communication (LOCC) measurement scheme. Two types of identification problems are considered when a single copy of each reference state is available. We show that a LOCC scheme attains the globally achievable identification probability in the minimum-error identification problem. In the unambiguous identification problem, however, the maximal success probability by means of LOCC is shown to be less than the globally achievable identification probability.

I. INTRODUCTION

It is an extremely nontrivial problem to distinguish different states of a quantum system by measurement [1, 2, 3, 4, 5]. First of all, this is because of the statistical nature of quantum measurement, which destroys the state of the system to be measured and one cannot clone an unknown quantum state [6]. Another relevant issue is nonlocality of quantum mechanics. When the system to be measured is a composite, we can generally obtain more information about the system by a global measurement on the whole system than by a combination of local measurements on its subsystems [7, 8, 9, 10, 11, 12].

Let us focus on the problem of distinguishing two pure states of a composite system which is shared by two parties. It is a fundamental question of quantum-information theory whether the optimal discrimination can be obtained by means of a local operations and classical communication (LOCC) scheme for the two parties.

Walgate *et al.* [13] showed that any two mutually orthogonal pure states can be perfectly distinguished by LOCC. This is rather surprising since their result holds regardless of entanglement of the states. It has also been shown that any two generally nonorthogonal pure states can be optimally discriminated by LOCC: the optimal success probability of discrimination by a global measurement can be attained by a LOCC protocol. This was shown for two types of discrimination problems: the discrimination problem with minimum error [14] where an erroneous guess is allowed, and the unambiguous discrimination problem [15, 16, 17] where no error is allowed but an inconclusive result can be produced. These results can be interpreted as meaning that there is no nonlocality in the discrimination of two pure states.

We can consider a different setting for the discrimination problem of two pure states. In the usual setting, it is assumed that perfect classical knowledge of the two states to be discriminated (ρ_1 and ρ_2) is given. The measurement scheme for the optimal discrimination naturally

depends on the classical knowledge of the states. Instead, let us assume that there is no classical knowledge of the states ρ_1 and ρ_2 , but a certain number (N) of their copies are available as reference states. One's task is to correctly identify a given input state ρ with one of the reference states ρ_1 and ρ_2 by means of a measurement on the whole state $\rho \otimes \rho_1^{\otimes N} \otimes \rho_2^{\otimes N}$. When the number of copies N is infinite, the problem is reduced to quantum state discrimination. This is because we can always obtain complete classical knowledge of a quantum state if we have infinitely many copies of the state. We call this problem "quantum state identification." The optimal success probability by a global measurement scheme has been determined for the minimum-error [18] and unambiguous [19, 20] identification problems.

In this paper, we investigate the pure-state identification problem of $N = 1$ where the two reference pure states ρ_1 and ρ_2 are bipartite. The input state ρ given to Alice and Bob is guaranteed to be one of the reference states ρ_1 and ρ_2 with prior probabilities η_1 and η_2 . The two reference states are assumed to be independently and randomly distributed on the pure-state space in a unitary invariant way. Each reference state generated in this way is generally entangled. Two types of identification problems will be considered: identification with minimum error and unambiguous identification where no error is allowed. For the minimum-error identification, we will demonstrate that Alice and Bob can identify the input state by means of a LOCC protocol with the success probability given by the optimal global identification scheme. In the case of the unambiguous identification problem, the maximal success probability by means of LOCC will be shown to be less than the globally achievable identification probability.

II. MINIMUM-ERROR IDENTIFICATION WITHOUT LOCC CONDITIONS

In this section, we will precisely formulate the minimum-error pure-state identification problem and derive the maximum mean success probability without the LOCC conditions for the case of $N = 1$ and an arbitrary prior occurrence probability of the reference states. In the case of a single-qubit system, the problem has been solved by Bergou *et al.* [21]. For the case of general N but with equal prior occurrence probabilities, see Ref. [18].

We have three quantum systems numbered 0, 1, and 2, each on a d -dimensional space \mathbf{C}^d . The input pure state $\rho = |\phi\rangle\langle\phi|$ is prepared in system 0 and the two reference pure states $\rho_1 = |\phi_1\rangle\langle\phi_1|$ and $\rho_2 = |\phi_2\rangle\langle\phi_2|$ in systems 1 and 2, respectively. The space that an operator acts on is specified by a number in parentheses. For example, $\rho(0)$ is a density operator on system 0. The input state ρ is promised to be one of the reference states ρ_1 and ρ_2 with prior probabilities $\{\eta_1, \eta_2\}$. The two reference states are independently and randomly chosen from the state space \mathbf{C}^d in a unitary invariant way. More precisely, the distribution is assumed uniform on the $(2d - 1)$ -dimensional unit hypersphere of $2d$ real variables $\{\text{Re}c_i, \text{Im}c_i\}_{i=0}^{d-1}$, where c_i are the expansion coefficients of the state in terms of an orthonormal base $\{|i\rangle\}_{i=0}^{d-1}$. The distribution does not depend on a particular choice of the base.

Our task is to correctly identify the input state with one of the reference states $\rho_\mu (\mu = 1, 2)$ by measuring the whole system $0 \otimes 1 \otimes 2$. We denote the corresponding positive-operator-valued measure (POVM) elements by $E_\mu (\mu = 1, 2)$. The mean identification success probability is then given by

$$p(d) = \sum_{\mu=1,2} \eta_\mu \left\langle \text{tr}[E_\mu \rho_\mu(0) \rho_1(1) \rho_2(2)] \right\rangle, \quad (1)$$

where the symbol $\langle \dots \rangle$ represents the average over the reference states ρ_1 and ρ_2 . Note that the POVM E_μ is independent of ρ_1 and ρ_2 , since we have no classical knowledge of the reference states.

The average over the reference states can be readily performed by using the formula [22]:

$$\langle \rho^{\otimes n} \rangle = \frac{\mathcal{S}_n}{d_n}, \quad (2)$$

where \mathcal{S}_n is the projector on to the totally symmetric subspace of $(\mathbf{C}^d)^{\otimes n}$ and d_n is its dimension given by $d_n = {}_{n+d-1}C_{d-1}$. Using $E_2 = 1 - E_1$, the mean success probability to be maximized is written as

$$p(d) = \eta_2 + \frac{1}{d_1 d_2} \text{tr}[E_1 (\eta_1 \mathcal{S}(01) - \eta_2 \mathcal{S}(02))], \quad (3)$$

where $\mathcal{S}(01)$ and $\mathcal{S}(02)$ are the projectors onto the totally symmetric subspaces of spaces $0 \otimes 1$ and $0 \otimes 2$, respectively.

The only restriction on the POVM element E_1 is $0 \leq E_1 \leq 1$. In order to maximize the mean success probability Eq.(3), we use the following result which holds for any Hermitian operator Δ :

$$\begin{aligned} & \max_{0 \leq E \leq 1} \text{tr}[E\Delta] \\ &= \text{sum of all positive eigenvalues of } \Delta, \end{aligned}$$

where the maximum is attained when E is the projector P_+ onto the subspace $V_+(\Delta)$ spanned by all eigenstates of Δ with a positive eigenvalue. Note that it does not matter whether the subspace $V_+(\Delta)$ includes eigenvectors with zero eigenvalue. In our case, Δ is defined to be

$$\Delta = \eta_1 \mathcal{S}(01) - \eta_2 \mathcal{S}(02). \quad (4)$$

Let us decompose the total space into three subspaces according to the symmetry with respect to system permutations [23]:

$$V = \mathbf{C}^d \otimes \mathbf{C}^d \otimes \mathbf{C}^d = V_S \oplus V_A \oplus V_M. \quad (5)$$

Here V_S is the totally symmetric subspace of dimension $\dim V_S \equiv d_3 = d(d+1)(d+2)/6$ and V_A is the totally antisymmetric subspace of dimension $\dim V_A = d(d-1)(d-2)/6$. The remaining subspace V_M is the mixed symmetric subspace of dimension $\dim V_M = 2d(d^2 - 1)/3$. The subspace V_M contains the two-dimensional irreducible representation of the symmetric group of order 3, S_3 , with multiplicity $\dim V_M/2$. We denote projectors onto V_S , V_A , and V_M by \mathcal{S}_3 , \mathcal{A}_3 and \mathcal{M}_3 , respectively.

It is clear that $\Delta = \eta_1 - \eta_2$ in V_S and $\Delta = 0$ in V_A . To determine eigenvalues of Δ in V_M , it is convenient to introduce two operators D and A :

$$D \equiv \mathcal{S}(01) - \mathcal{S}(02) = \frac{1}{2} (T(01) - T(02)), \quad (6)$$

$$A \equiv \mathcal{S}(01) + \mathcal{S}(02) - 1 = \frac{1}{2} (T(01) + T(02)). \quad (7)$$

Here, $T(01)$ is the operator that exchanges systems 0 and 1 and $T(02)$ exchanges systems 0 and 2. Calculating D^2 , we find

$$\begin{aligned} D^2 &= \frac{1}{4} (2 - T(01)T(02) - T(02)T(01)) \\ &= \frac{3}{4} (1 - \mathcal{S}_3 - \mathcal{A}_3) \\ &= \frac{3}{4} \mathcal{M}_3, \end{aligned}$$

which implies that eigenvalues of D are $\pm\sqrt{3}/2$ in V_M and 0 otherwise. It is also easy to show that

$$DA + AD = 0, \quad (8)$$

$$A^2 = 1 - D^2. \quad (9)$$

The anticommutability of Eq.(8) implies that if $|+\rangle$ is an eigenstate of D with eigenvalue $\sqrt{3}/2$, then $A|+\rangle$ is also

an eigenstate of D with eigenvalue $-\sqrt{3}/2$. By Eq.(9), we find that $|- \rangle \equiv 2A|+ \rangle$ is correctly normalized. Note that the positive and negative eigenvalues of D have the same multiplicity. Thus we can choose the orthonormal base $\{|+, k \rangle, |-, k \rangle\}$ in $V_{\mathcal{M}}$ such that

$$\begin{aligned} D|+, k \rangle &= +\frac{\sqrt{3}}{2}|+, k \rangle, \\ D|-, k \rangle &= -\frac{\sqrt{3}}{2}|-, k \rangle, \\ A|+, k \rangle &= \frac{1}{2}|-, k \rangle, \\ A|-, k \rangle &= \frac{1}{2}|+, k \rangle, \end{aligned}$$

where the index k runs from 1 to $\dim V_{\mathcal{M}}/2$. In this base, D and A are block-diagonalized with respect to k and each block has the following 2×2 matrix representation:

$$D = \begin{pmatrix} \frac{\sqrt{3}}{2} & 0 \\ 0 & -\frac{\sqrt{3}}{2} \end{pmatrix}, \quad A = \begin{pmatrix} 0 & \frac{1}{2} \\ \frac{1}{2} & 0 \end{pmatrix}.$$

In terms of D and A , the operator Δ is written as

$$\Delta = \frac{1}{2}(\eta_1 - \eta_2 + D + (\eta_1 - \eta_2)A). \quad (10)$$

The operator Δ is also block-diagonalized with the same 2×2 matrix representation which can be readily diagonalized. Two eigenvalues of Δ are given by

$$\lambda_{\pm} = \frac{1}{2}(\eta_1 - \eta_2 \pm \sqrt{1 - \eta_1\eta_2}), \quad (11)$$

and we find that $\lambda_+ \geq 0$ and $\lambda_- \leq 0$.

Now we can calculate the maximum success probability. Let us assume $\eta_1 \geq \eta_2$ for the moment. The positive eigenvalues of Δ are $\eta_1 - \eta_2$ in V_S with multiplicity $\dim V_S$ and λ_+ in $V_{\mathcal{M}}$ with multiplicity $\dim V_{\mathcal{M}}/2$. We thus obtain

$$\begin{aligned} p_{\max}(d) &= \eta_2 + \frac{1}{d_1 d_2} \left((\eta_1 - \eta_2) \dim V_S + \lambda_+ \frac{\dim V_{\mathcal{M}}}{2} \right) \\ &= \frac{1}{2} + \frac{d+2}{6d}(\eta_1 - \eta_2) + \frac{d-1}{3d}\sqrt{1 - \eta_1\eta_2}. \end{aligned} \quad (12)$$

If $\eta_1 \leq \eta_2$, the only positive eigenvalue of Δ is λ_+ in $V_{\mathcal{M}}$; hence we obtain

$$\begin{aligned} p_{\max}(d) &= \eta_2 + \frac{1}{d_1 d_2} \lambda_+ \frac{\dim V_{\mathcal{M}}}{2} \\ &= \frac{1}{2} - \frac{d+2}{6d}(\eta_1 - \eta_2) + \frac{d-1}{3d}\sqrt{1 - \eta_1\eta_2}. \end{aligned} \quad (13)$$

These two cases can be combined to yield a symmetric form of the maximum success identification probability for the general magnitude relation between η_1 and η_2 :

$$p_{\max}(d) = \frac{1}{2} + \frac{d+2}{6d}|\eta_1 - \eta_2| + \frac{d-1}{3d}\sqrt{1 - \eta_1\eta_2}. \quad (14)$$

The maximum is attained when the POVM element E_1 is given by P_+ , the projector onto the subspace of positive eigenvalues of Δ . The $p_{\max}(d)$ given by Eq.(14) reproduces the result for the case $d = 2$ obtained in Ref.[21] and the one for arbitrary d in Ref.[18] when $\eta_1 = \eta_2 = 1/2$.

III. MINIMUM-ERROR IDENTIFICATION BY LOCC

Let us assume that each of the three systems 0, 1, and 2, where the input state and the two reference states are prepared, consists of two subsystems. The state space of each system is represented by a tensor product $\mathbf{C}^d = \mathbf{C}^{d_a} \otimes \mathbf{C}^{d_b}$, which is shared by Alice and Bob. Their task is to identify a given input bipartite state with one of the two bipartite reference states by means of local operations and classical communication. As in the preceding section, the two reference states are chosen independently and randomly from the pure state space \mathbf{C}^d in a unitary invariant way. Therefore, those bipartite states are generally entangled. The question is whether Alice and Bob can achieve the maximum mean identification success probability given by the global measurement scheme. In this section, we will show that this is possible by explicitly constructing a LOCC protocol which achieves it.

The mean success probability is given by Eq.(3) in the preceding section. The optimal global POVM element E_1 is P_+ , the projector onto the subspace of positive eigenvalues of Δ defined by Eq.(4). The projector P_+ does not apparently satisfy the conditions of LOCC, since the operator Δ is not of a separable form. However, it should be noticed that $\text{tr}[E_1 \Delta]$ remains the same if the support of E_1 contains states with zero eigenvalue of Δ . It is this freedom that we will exploit in order to construct a POVM element E_1 that satisfies the LOCC conditions.

We begin by rewriting the operator Δ of Eq.(10) in terms of local operators of Alice and Bob. Note that the exchange operator $T(01)$, for example, can be written as $T(01) = T^{(a)}(01) \otimes T^{(b)}(01)$, where $T^{(a)}(01)$ is the operator that exchanges Alice's part of systems 0 and 1 and $T^{(b)}(01)$ is defined for Bob's part in the same way. Hereafter, we use the superscript (a) or (b) for an operator to indicate which space (of Alice or Bob) the operator acts on. Since we have

$$D = D^{(a)} \otimes A^{(b)} + A^{(a)} \otimes D^{(b)}, \quad (15)$$

$$A = D^{(a)} \otimes D^{(b)} + A^{(a)} \otimes A^{(b)}, \quad (16)$$

the operator Δ is expressed as

$$\begin{aligned} \Delta &= \frac{1}{2} \left(\eta_1 - \eta_2 + D^{(a)} A^{(b)} + A^{(a)} D^{(b)} \right. \\ &\quad \left. + (\eta_1 - \eta_2)(D^{(a)} D^{(b)} + A^{(a)} A^{(b)}) \right). \end{aligned} \quad (17)$$

The task for Alice and Bob is to maximize $\text{tr}[E_1^L \Delta]$ with a POVM element E_1^L which satisfies LOCC condi-

tions. We first construct a separable POVM E_1^L which attains the maximum value $\text{tr}[P_+ \Delta]$. This separable POVM E_1^L will then be shown to satisfy the LOCC conditions. Without loss of generality, we assume $\eta_1 \leq \eta_2$ throughout this section, since the problem is symmetric with respect to ρ_1 and ρ_2 .

Suppose that Alice and Bob first determine the permutation symmetry of their systems by a projective measurement with projection operators $\{\mathcal{S}_3^{(a)}, \mathcal{A}_3^{(a)}, \mathcal{M}_3^{(a)}\}$ and $\{\mathcal{S}_3^{(b)}, \mathcal{A}_3^{(b)}, \mathcal{M}_3^{(b)}\}$, respectively. If one of them found that his or her system is totally symmetric or antisymmetric, it is easy for the other party to find the best strategy. For example, assume that Alice found her system to be totally symmetric. Knowing Alice's outcome, Bob performs a POVM measurement, which we denote by $x^{(b)}$. The contribution to $\text{tr}[E_1^L \Delta]$ is then given by

$$\text{tr}[\mathcal{S}_3^{(a)} \otimes x^{(b)} \Delta] = \dim(V_S^{(a)}) \text{tr}_b[x^{(b)} \Delta^{(b)}],$$

since $\mathcal{S}_3^{(a)} D^{(a)} = 0$ and $\mathcal{S}_3^{(a)} A^{(a)} = \mathcal{S}_3^{(a)}$. It is clear that the best strategy for Bob is to take the projector $P_+^{(b)}$ onto the positive-eigenvalue space of $\Delta^{(b)}$. Note that the positive-eigenvalue space of $\Delta^{(b)}$ is a subspace of $V_M^{(b)}$, since the eigenvalue of $\Delta^{(b)}$ in $V_S^{(b)}$ is $\eta_1 - \eta_2 (\leq 0)$. In this case the contribution to $\text{tr}[E_1^L \Delta]$ is given by

$$\text{tr}[\mathcal{S}_3^{(a)} \otimes P_+^{(b)} \Delta] = \frac{\lambda_+}{2} \dim V_S^{(a)} \dim V_M^{(b)}.$$

If Alice's part is totally antisymmetric, the operator for Bob is given by

$$\begin{aligned} \text{tr}_a[\mathcal{A}_3^{(a)} \Delta] &= \dim(V_A^{(a)}) \Delta'^{(b)}, \\ \Delta'^{(b)} &\equiv \frac{1}{2} (\eta_1 - \eta_2 - D^{(b)} - (\eta_1 - \eta_2) A^{(b)}). \end{aligned}$$

The operator $\Delta'^{(b)}$ differs from $\Delta^{(b)}$ only in the signs in front of $D^{(b)}$ and $A^{(b)}$. Its eigenvalues are 0 in $V_S^{(b)}$ and $\eta_1 - \eta_2 (\leq 0)$ in $V_A^{(b)}$. In $V_M^{(b)}$, the operator $\Delta'^{(b)}$ has eigenvalue λ_- in the positive-eigenvalue subspace of $\Delta^{(b)}$ and λ_+ in the negative-eigenvalue subspace of $\Delta^{(b)}$. This implies that Bob's best POVM element is $P_-^{(b)}$, the projector onto the $\Delta^{(b)}$'s negative-eigenvalue subspace in $V_M^{(b)}$. The contribution to $\text{tr}[E_1^L \Delta]$ in this case is given by

$$\text{tr}[\mathcal{A}_3^{(a)} \otimes P_-^{(b)} \Delta] = \frac{\lambda_+}{2} \dim V_A^{(a)} \dim V_M^{(b)}.$$

The same argument also holds when Bob's system is totally symmetric or antisymmetric. Therefore, when the total state does not belong to $V_M^{(a)} \otimes V_M^{(b)}$, the whole contribution to $\text{tr}[E_1^L \Delta]$ is given by

$$\begin{aligned} &\text{tr}\left[\left(\mathcal{S}_3^{(a)} P_+^{(b)} + \mathcal{A}_3^{(a)} P_-^{(b)} + P_+^{(a)} \mathcal{S}_3^{(b)} + P_-^{(a)} \mathcal{A}_3^{(b)}\right) \Delta\right] \\ &= \frac{1}{2} \lambda_+ \left(\dim V_S^{(a)} \dim V_M^{(b)} + \dim V_A^{(a)} \dim V_M^{(b)} \right. \\ &\quad \left. + \dim V_M^{(a)} \dim V_S^{(b)} + \dim V_M^{(a)} \dim V_A^{(b)} \right). \end{aligned} \quad (18)$$

When the total state belongs to $V_M^{(a)} \otimes V_M^{(b)}$, construction of the best strategy for Alice and Bob is rather involved. First we introduce the following operators X_1 and X_2 for each of Alice's and Bob's spaces:

$$\begin{aligned} X_1^{(\kappa)} &= \frac{2}{\sqrt{3}} D^{(\kappa)}, \\ X_2^{(\kappa)} &= 2A^{(\kappa)}, \quad (\kappa = a, b). \end{aligned} \quad (19)$$

Note that $X_1^{(\kappa)}$ and $X_2^{(\kappa)}$ anticommute and $(X_1^{(\kappa)})^2 = (X_2^{(\kappa)})^2 = 1$ in the mixed symmetric space $V_M^{(\kappa)}$. The operator Δ in terms of $X_i^{(\kappa)}$ is not diagonal with respect to the index i . We further define rotated $X_i^{(\kappa)}$'s in order to diagonalize Δ with respect to the index i :

$$\begin{aligned} Y_1^{(\kappa)} &= \cos \theta X_1^{(\kappa)} + \sin \theta X_2^{(\kappa)}, \\ Y_2^{(\kappa)} &= -\sin \theta X_1^{(\kappa)} + \cos \theta X_2^{(\kappa)}, \quad (\kappa = a, b). \end{aligned} \quad (20)$$

We find that Δ takes the following "diagonal" form:

$$\Delta = \frac{1}{2} \left(\eta_1 - \eta_2 + \lambda_+ Y_1^{(a)} Y_1^{(b)} + \lambda_- Y_2^{(a)} Y_2^{(b)} \right), \quad (21)$$

if we take

$$\begin{aligned} \cos 2\theta &= \frac{\eta_1 - \eta_2}{2\sqrt{1 - \eta_1 \eta_2}}, \\ \sin 2\theta &= \frac{\sqrt{3}}{2\sqrt{1 - \eta_1 \eta_2}}. \end{aligned}$$

The eigenvalues of $Y_i^{(\kappa)}$ are 1 and -1 with multiplicity $\dim V_M^{(\kappa)}/2$ since we have

$$\begin{aligned} (Y_1^{(\kappa)})^2 &= 1, \quad (Y_2^{(\kappa)})^2 = 1, \\ Y_1^{(\kappa)} Y_2^{(\kappa)} + Y_2^{(\kappa)} Y_1^{(\kappa)} &= 0. \end{aligned} \quad (22)$$

The positive- and negative-eigenvalue subspaces of $Y_1^{(\kappa)}$ are transformed into each other by the operation of $Y_2^{(\kappa)}$ and vice versa. We should also notice that $|\lambda_-| \geq |\lambda_+|$ when $\eta_1 \leq \eta_2$. These considerations imply that the optimal separate POVM element is given by $Q_+^{(a)} \otimes Q_-^{(b)} + Q_-^{(a)} \otimes Q_+^{(b)}$, where $Q_\pm^{(\kappa)}$ is the projector onto the positive- and negative-eigenvalue subspaces of $Y_2^{(\kappa)}$. The contribution to $\text{tr}[E_1^L \Delta]$ is found to be

$$\begin{aligned} &\text{tr}\left[\left(Q_+^{(a)} \otimes Q_-^{(b)} + Q_-^{(a)} \otimes Q_+^{(b)}\right) \Delta\right] \\ &= \frac{1}{4} (\eta_1 - \eta_2 - \lambda_-) \dim V_M^{(a)} \dim V_M^{(b)} \\ &= \frac{1}{4} \lambda_+ \dim V_M^{(a)} \dim V_M^{(b)}, \end{aligned} \quad (23)$$

where we used $\text{tr}[Q_\pm^{(\kappa)} Y_1^{(\kappa)}] = 0$.

Thus the whole POVM element is given by

$$\begin{aligned} E_1^L &= \mathcal{S}_3^{(a)} P_+^{(b)} + \mathcal{A}_3^{(a)} P_-^{(b)} + P_+^{(a)} \mathcal{S}_3^{(b)} + P_-^{(a)} \mathcal{A}_3^{(b)} \\ &\quad + Q_+^{(a)} Q_-^{(b)} + Q_-^{(a)} Q_+^{(b)}. \end{aligned} \quad (24)$$

Adding Eqs.(18) and (23), we find that $\text{tr}[E_1^L \Delta]$ indeed attains the maximum value given by the global POVM element $E_1 = P_+$:

$$\text{tr}[E_1^L \Delta] = \frac{1}{2} \lambda_+ \dim V_{\mathcal{M}} = \text{tr}[P_+ \Delta]. \quad (25)$$

To show the above equality, we used the relation

$$\begin{aligned} \dim V_{\mathcal{M}} &= \dim V_S^{(a)} \dim V_{\mathcal{M}}^{(b)} + \dim V_{\mathcal{M}}^{(a)} \dim V_S^{(b)} \\ &\quad + \dim V_A^{(a)} \dim V_{\mathcal{M}}^{(b)} + \dim V_{\mathcal{M}}^{(a)} \dim V_A^{(b)} \\ &\quad + \frac{1}{2} \dim V_{\mathcal{M}}^{(a)} \dim V_{\mathcal{M}}^{(b)}, \end{aligned} \quad (26)$$

which can be readily verified by a straightforward calculation. This relation can be also understood from the viewpoint of inner (Kronecker) products of two representations of the symmetric group of order 3: the product of two mixed symmetric representations contains the totally symmetric and antisymmetric representations in addition to the mixed symmetric representation, whereas the product of the totally (anti)symmetric representation and the mixed symmetric representation is the mixed symmetric representation.

We can show that the POVM element E_1^L given in Eq.(24) can be implemented with an LOCC protocol. First Alice and Bob determine which permutation symmetries each one's local state has; totally symmetric, totally antisymmetric, or mixed symmetric. If one of them finds that his or her state is totally symmetric or antisymmetric and the other party's state is mixed symmetric, this party with the mixed symmetric state performs the measurement by the projectors $\{P_+^{(\kappa)}, P_-^{(\kappa)}\}$. They conclude that the input state is ρ_1 if the combination of their outcomes is either (symmetric, P_+) or (antisymmetric, P_-). Otherwise they conclude that the input state is ρ_2 . When Alice and Bob find both the local states are mixed symmetric, they perform the projection measurement by $\{Q_+^{(\kappa)}, Q_-^{(\kappa)}\}$. They conclude that the input state is ρ_1 , only when the combination of outcomes is "(+, -)".

Thus we conclude that minimum-error pure-state identification with any prior occurrence probability can be optimally performed within the LOCC scheme.

IV. UNAMBIGUOUS IDENTIFICATION AND SYMMETRIES OF POVM

In this section, we precisely formulate the unambiguous identification problem of two pure states, and re-derive the optimal success probability attainable by the global measurement scheme, though the result has been reported in Refs.[19, 20]. The purpose of this section is to explain two important symmetries of the measurement scheme of this problem, which will also play a crucial role in determining the optimal probability by the LOCC scheme in the next section. In this and the next sections, we assume equal prior probabilities for the two reference states.

A. Problem

In the unambiguous identification problem, the task is to unambiguously identify the input pure state ρ with one of the two pure reference states ρ_1 and ρ_2 . Our measurement can produce three outcomes $\mu = 1, 2, 0$. If the outcome is $\mu (= 1, 2)$, we are certain that the input state ρ is ρ_μ , and outcome 0 means that we are not certain about the identity of the input; this is called an inconclusive result. Let us introduce a POVM, $\{E_\mu\}_{\mu=0,1,2}$ corresponding to the three measurement outcomes. The mean success probability of identification is then given by

$$p = \frac{1}{2} \left\langle \text{tr}[E_1 \rho_1(0) \rho_1(1) \rho_2(2)] + \text{tr}[E_2 \rho_2(0) \rho_1(1) \rho_2(2)] \right\rangle. \quad (27)$$

The condition that we are not allowed to make a mistake imposes the following no-error conditions on E_1 and E_2 :

$$\begin{cases} \text{tr}[E_1 \rho_2(0) \rho_1(1) \rho_2(2)] = 0, \\ \text{tr}[E_2 \rho_1(0) \rho_1(1) \rho_2(2)] = 0, \end{cases} \quad \text{for any } \rho_1 \text{ and } \rho_2. \quad (28)$$

In what follows, we will optimize the mean success probability of Eq.(27) under the no-error conditions Eq.(28). Performing the average over the reference states by Eq.(2), we write the mean success probability (27) as

$$p(d) = \frac{1}{2d_2 d_1} (\text{tr}[E_1 \mathcal{S}(01)] + \text{tr}[E_2 \mathcal{S}(02)]), \quad (29)$$

where $\mathcal{S}(01)$ and $\mathcal{S}(02)$ are the projectors onto the totally symmetric subspaces of spaces $0 \otimes 1$ and $0 \otimes 2$, respectively. Averaging the no-error conditions of Eq.(28), we obtain

$$\text{tr}[E_1 \mathcal{S}(02)] = 0, \quad \text{tr}[E_2 \mathcal{S}(01)] = 0. \quad (30)$$

Since E_1 and $\mathcal{S}(02)$ are both positive operators, the above conditions imply that the supports of them are orthogonal to each other. The same is true for E_2 and $\mathcal{S}(01)$. The no-error conditions are thus equivalent to

$$E_1 \mathcal{S}(02) = \mathcal{S}(02) E_1 = 0, \quad (31)$$

$$E_2 \mathcal{S}(01) = \mathcal{S}(01) E_2 = 0. \quad (32)$$

B. Symmetries of POVM

The set of POVMs satisfying the no-error conditions is convex; if two POVMs E_μ and E'_μ respect the no-error conditions, so does their convex linear combination $rE_\mu + (1-r)E'_\mu$ for any $0 \leq r \leq 1$. The resulting success probability is also a convex combination, $p(rE + (1-r)E') = rp(E) + (1-r)p(E')$, with an obviously abbreviated notation. It is this convexity of the POVM that we exploit in order to impose two symmetries on the optimal POVM without loss of generality.

First we consider the exchange symmetry between systems 1 and 2. For an optimal POVM F_μ , we define another POVM by

$$\begin{aligned} F'_1 &= T(12)F_2T(12), \\ F'_2 &= T(12)F_1T(12), \\ F'_0 &= T(12)F_0T(12), \end{aligned}$$

where $T(12)$ is the exchange operator between systems 1 and 2. The POVM F'_μ is clearly legitimate and optimal. Furthermore, a new POVM $E_\mu = \frac{1}{2}(F_\mu + F'_\mu)$, which is a convex linear combination of F_μ and F'_μ , is also optimal and satisfies the exchange symmetry between systems 1 and 2:

$$\begin{aligned} E_1 &= T(12)E_2T(12), \\ E_2 &= T(12)E_1T(12), \\ E_0 &= T(12)E_0T(12). \end{aligned} \quad (33)$$

The second important symmetry is the unitary symmetry of the distribution of the reference states. If a POVM F_μ is optimal, another POVM defined by

$$U^{\otimes 3}F_\mu(U^{\otimes 3})^{-1},$$

is also legitimate and optimal for arbitrary unitary operator U . We now construct a POVM by

$$E_\mu = \int dU U^{\otimes 3} F_\mu(U^{\otimes 3})^{-1}, \quad (34)$$

where dU is the normalized positive invariant measure of the group $U(d)$. The POVM E_μ is clearly a legitimate and optimal POVM. We can show that E_μ commutes with $U^{\otimes 3}$ for any unitary U :

$$\begin{aligned} U^{\otimes 3}E_\mu &= \int dU' (UU')^{\otimes 3} F_\mu(UU'^{\otimes 3})^{-1} U^{\otimes 3} \\ &= \int dU' U'^{\otimes 3} F_\mu(U'^{\otimes 3})^{-1} U^{\otimes 3} \\ &= E_\mu U^{\otimes 3}, \end{aligned} \quad (35)$$

which means that E_μ is a scalar with respect to the group $U(d)$. Thus we can assume that the optimal POVM satisfies the exchange symmetry of Eq.(33) and is scalar with respect to the group $U(d)$.

By the exchange symmetry, the mean success probability to be optimized takes the form

$$p = \frac{1}{d_2 d_1} \text{tr}[E_1 \mathcal{S}(01)], \quad (36)$$

where E_1 is a unitary scalar and subject to the conditions

$$E_1 \geq 0, \quad 1 \geq E_1 + T(12)E_1T(12), \quad (37)$$

in addition to the no-error conditions given by Eq.(31).

C. Optimal identification probability

As in the two preceding sections, we decompose the total space into three subspaces V_S , V_A , and V_M , according to the symmetry with respect to system permutations. Under the unitary transformation $U^{\otimes 3}$, the three subspaces V_S , V_A , and V_M are clearly invariant since $U^{\otimes 3}$ commutes with system permutations. Furthermore, it is known that $U^{\otimes 3}$ acts on V_S and V_A irreducibly, and the mixed symmetric space V_M contains two $(\dim V_M/2)$ -dimensional irreducible representations of the group $U(d)$ [23]. Now, suppose that a positive operator E is a unitary scalar; E commutes with $U^{\otimes 3}$. If $\text{tr}[E\mathcal{S}(02)] = 0$, Schur's lemma requires that E be a linear combination of two projection operators \mathcal{A}_3 and $\mathcal{M}_3\mathcal{A}(02)$, with $\mathcal{A}(02) = 1 - \mathcal{S}(02)$ being the projector onto the antisymmetric subspace of $0 \otimes 2$. Similarly, if $\text{tr}[E\mathcal{A}(02)] = 0$, E is a linear combination of \mathcal{S}_3 and $\mathcal{M}_3\mathcal{S}(02)$. These facts will be used also in the next section.

With these considerations, we can determine the operator form of the POVM element E_1 . The operator E_1 is given by a linear combination of \mathcal{A}_3 and $\mathcal{M}_3\mathcal{A}(02)$ since E_1 is a unitary scalar and satisfies the no-error condition $\text{tr}[E_1\mathcal{S}(02)] = 0$. But \mathcal{A}_3 does not contribute to the mean success probability of Eq.(36). Thus, without loss of generality, we can write

$$E_1 = \alpha \mathcal{M}_3\mathcal{A}(02), \quad (38)$$

where α is a positive coefficient. The range of α is restricted by the positivity of E_0 , which is written as

$$\begin{aligned} E_1 + T(12)E_1T(12) &= \alpha \mathcal{M}_3(\mathcal{A}(02) + \mathcal{A}(01)) \\ &= \alpha \mathcal{M}_3(1 - A) \leq 1. \end{aligned}$$

This requires $\alpha \leq 2/3$, since the operator A defined in Eq.(7) has eigenvalues $\pm 1/2$ in V_M . Clearly, the mean success probability attains its maximum when α takes the largest possible value, $2/3$. The optimal POVM is thus given by

$$\begin{aligned} E_1 &= \frac{2}{3} \mathcal{M}_3\mathcal{A}(02), \\ E_2 &= \frac{2}{3} \mathcal{M}_3\mathcal{A}(01), \\ E_0 &= \frac{1}{3} \mathcal{M}_3(1 + 2A) + \mathcal{S}_3 + \mathcal{A}_3. \end{aligned} \quad (39)$$

In order to obtain the optimal probability, we need trace $\text{tr}[\mathcal{M}_3\mathcal{A}(02)\mathcal{S}(01)]$, which is calculated as follows:

$$\begin{aligned} &\text{tr}[\mathcal{M}_3\mathcal{A}(02)\mathcal{S}(01)] \\ &= \frac{1}{4} \text{tr}[\mathcal{M}_3(1 + T(01) - T(02) - T(02)T(01))] \\ &= \frac{1}{2} \text{tr}[\mathcal{M}_3 D^2] = \frac{3}{8} \dim V_M, \end{aligned} \quad (40)$$

where we used $\text{tr}[\mathcal{M}_3T(01)] = \text{tr}[\mathcal{M}_3T(02)] = 0$. Using the explicit expressions for the dimensions, we finally

obtain the optimal mean success probability of unambiguous identification:

$$p_{\max} = \frac{d-1}{3d}. \quad (41)$$

V. LOCAL UNAMBIGUOUS IDENTIFICATION

Let us now assume that each of the three systems consists of two subsystems shared by Alice and Bob and its state space is represented by a tensor product $\mathcal{C}^d = \mathcal{C}^{d_a} \otimes \mathcal{C}^{d_b}$. The task of Alice and Bob is to unambiguously identify a given input state by means of LOCC with one of the two reference states.

A. Separable POVM and symmetries

Any POVM E_μ^L which satisfies the LOCC conditions has a separable form:

$$E_\mu^L = \sum_i E_{\mu i}^{(a)} \otimes E_{\mu i}^{(b)}, \quad (\mu = 0, 1, 2), \quad (42)$$

where

$$E_{\mu i}^{(a)} \geq 0, \quad E_{\mu i}^{(b)} \geq 0, \quad \sum_\mu E_\mu^L = 1. \quad (43)$$

It is known that there exist separable POVMs that do not satisfy the LOCC conditions [11]. We will first optimize the success probability within the separable class of POVMs, and then show that the obtained optimal separable POVM can be implemented by a LOCC protocol.

Note that a convex linear combination of separable POVMs is again separable, and the no-error conditions $\text{tr}[E_1^L \mathcal{S}(02)] = \text{tr}[E_2^L \mathcal{S}(01)] = 0$ are also preserved. This enables us to impose two symmetries on the optimal separable POVM as in the preceding section. We begin with the exchange symmetry for systems 1 and 2. Since $T(12) = T^{(a)}(12) \otimes T^{(b)}(12)$, it is clear that $T(12)E_\mu^L T(12)$ is separable if E_μ^L is separable:

$$\begin{aligned} & T(12)E_\mu^L T(12) \\ &= \sum_i T^{(a)}(12)E_{\mu i}^{(a)}T^{(a)}(12) \otimes T^{(b)}(12)E_{\mu i}^{(b)}T^{(b)}(12). \end{aligned}$$

Therefore, we can impose on separable POVMs the same exchange symmetry as the one given in Eq.(33):

$$\begin{aligned} E_1^L &= T(12)E_2^L T(12), \\ E_2^L &= T(12)E_1^L T(12), \\ E_0^L &= T(12)E_0^L T(12). \end{aligned} \quad (44)$$

For the unitary symmetry, we notice that $U^{\otimes 3}F_\mu^L(U^{\otimes 3})^{-1}$ is not generally separable for a separable POVM $F_\mu^L = \sum_i F_{\mu i}^{(a)} \otimes F_{\mu i}^{(b)}$. However, this

is true if U is a tensor product of two unitaries as $U = u^{(a)} \otimes v^{(b)}$:

$$\begin{aligned} & U^{\otimes 3}F_\mu^L(U^{\otimes 3})^{-1} \\ &= \sum_i u^{(a)\otimes 3}F_{\mu i}^{(a)}(u^{(a)\otimes 3})^{-1} \otimes v^{(b)\otimes 3}F_{\mu i}^{(b)}(v^{(b)\otimes 3})^{-1}. \end{aligned}$$

For the class of this separable U , we can repeat the argument given in the preceding section. Assume that a separable POVM F_μ^L is optimal. Integrating over $u^{(a)}$ and $v^{(b)}$ with the invariant measure, we obtain

$$\begin{aligned} E_\mu^L &\equiv \int du^{(a)} dv^{(b)} \left(u^{(a)\otimes 3} v^{(b)\otimes 3} \right) F_\mu^L \left(u^{(a)\otimes 3} v^{(b)\otimes 3} \right)^{-1} \\ &= \sum_i E_{\mu i}^{(a)} \otimes E_{\mu i}^{(b)}, \end{aligned}$$

where $E_{\mu i}^{(a)}$ and $E_{\mu i}^{(b)}$ are given by

$$\begin{aligned} E_{\mu i}^{(a)} &= \int du^{(a)} u^{(a)\otimes 3} F_{\mu i}^{(a)} (u^{(a)\otimes 3})^{-1}, \\ E_{\mu i}^{(b)} &= \int dv^{(b)} v^{(b)\otimes 3} F_{\mu i}^{(b)} (v^{(b)\otimes 3})^{-1}. \end{aligned}$$

The POVM E_μ^L obtained in this way is again separable and optimal. Furthermore, it is easy to see that $E_{\mu i}^{(a)}$ and $E_{\mu i}^{(b)}$ are both unitary scalar: for any unitaries $u^{(a)}$ and $v^{(b)}$, we have

$$[E_{\mu i}^{(a)}, u^{(a)}] = 0, \quad [E_{\mu i}^{(b)}, v^{(b)}] = 0. \quad (45)$$

Let us closely examine the no-error conditions for separable POVMs. The global projector $\mathcal{S}(02)$ is decomposed by local symmetry projectors as follows:

$$\mathcal{S}(02) = \mathcal{S}^{(a)}(02) \otimes \mathcal{S}^{(b)}(02) + \mathcal{A}^{(a)}(02) \otimes \mathcal{A}^{(b)}(02).$$

The no-error condition $\text{tr}[E_1^L \mathcal{S}(02)] = 0$ is then expressed as

$$\begin{aligned} & \sum_i \left(\text{tr}[E_{1i}^{(a)} \mathcal{S}^{(a)}(02)] \text{tr}[E_{1i}^{(b)} \mathcal{S}^{(b)}(02)] \right. \\ & \quad \left. + \text{tr}[E_{1i}^{(a)} \mathcal{A}^{(a)}(02)] \text{tr}[E_{1i}^{(b)} \mathcal{A}^{(b)}(02)] \right) = 0. \end{aligned}$$

In this equation, all terms are non-negative, implying that each term should vanish. Therefore, for each i , we have two possibilities: one is

$$\text{tr}[E_{1i}^{(a)} \mathcal{S}^{(a)}(02)] = 0, \text{ and } \text{tr}[E_{1i}^{(b)} \mathcal{A}^{(b)}(02)] = 0,$$

and the other is

$$\text{tr}[E_{1i}^{(a)} \mathcal{A}^{(a)}(02)] = 0, \text{ and } \text{tr}[E_{1i}^{(b)} \mathcal{S}^{(b)}(02)] = 0.$$

Note that other combinations like

$$\text{tr}[E_{1i}^{(a)} \mathcal{S}^{(a)}(02)] = 0, \text{ and } \text{tr}[E_{1i}^{(a)} \mathcal{A}^{(a)}(02)] = 0,$$

do not occur as this would imply that $E_{1i}^{(a)}$ or $E_{1i}^{(b)}$ is identically zero. From the no-error condition for E_2^L , we obtain the similar conditions for its components $E_{2i}^{(a)}$ and $E_{2i}^{(b)}$.

B. Possible operator form of separable POVM

As in the preceding section, we can show that a positive operator $E^{(p)}$ on space $V^{(p)}(p = a, b)$ which is a unitary scalar and satisfies $\text{tr}[E^{(p)}\mathcal{S}^{(p)}(02)] = 0$ is a linear combination of $\mathcal{A}_3^{(p)}$ and $\mathcal{M}_3^{(p)}\mathcal{A}^{(p)}(02)$. Similarly, if $E^{(p)}$ satisfies $\text{tr}[E^{(p)}\mathcal{A}^{(p)}(02)] = 0$, then $E^{(p)}$ can be written as a linear combination of $\mathcal{S}_3^{(p)}$ and $\mathcal{M}_3^{(p)}\mathcal{S}^{(p)}(02)$. Now we can write the possible form of the separable E_1^L which has the unitary symmetry and satisfies the no-error conditions:

$$\begin{aligned} E_1^L = & \alpha_1 \mathcal{S}_3^{(a)} \otimes \mathcal{M}_3^{(b)} \mathcal{A}^{(b)}(02) + \alpha_2 \mathcal{A}_3^{(a)} \otimes \mathcal{M}_3^{(b)} \mathcal{S}^{(b)}(02) \\ & + \alpha_3 \mathcal{M}_3^{(a)} \mathcal{S}^{(a)}(02) \otimes \mathcal{A}_3^{(b)} + \alpha_4 \mathcal{M}_3^{(a)} \mathcal{A}^{(a)}(02) \otimes \mathcal{S}_3^{(b)} \\ & + \beta_1 \mathcal{M}_3^{(a)} \mathcal{S}^{(a)}(02) \otimes \mathcal{M}_3^{(b)} \mathcal{A}^{(b)}(02) \\ & + \beta_2 \mathcal{M}_3^{(a)} \mathcal{A}^{(a)}(02) \otimes \mathcal{M}_3^{(b)} \mathcal{S}^{(b)}(02), \end{aligned} \quad (46)$$

where $\alpha_1, \alpha_2, \alpha_3, \alpha_4$ and β_1, β_2 are non-negative coefficients. The operators $\mathcal{S}_3 \otimes \mathcal{A}_3$ and $\mathcal{A}_3 \otimes \mathcal{S}_3$ are not included in E_1^L , since the corresponding outcomes do not arise; the total state contains no totally antisymmetric component. E_2^L and E_0^L are given by $E_2^L = T(12)E_1^L T(12)$ and $E_0^L = 1 - E_1^L - E_2^L$.

We have now found the possible operator form of a separable POVM that respects the no-error conditions. The remaining requirement on the POVM is the positivity of E_0^L , which restricts the range of the coefficients α and β . The positivity of E_0^L is equivalent to $E_1^L + E_2^L \leq 1$. We will separately check this inequality in each of all subspaces of the permutation symmetry.

In the subspace $V_S^{(a)} \otimes V_M^{(b)}$, the relevant part of $E_1^L + E_2^L$ is written as

$$\begin{aligned} & \alpha_1 (\mathcal{S}_3^{(a)} \otimes \mathcal{M}_3^{(b)} \mathcal{A}^{(b)}(02) + \mathcal{S}_3^{(a)} \otimes \mathcal{M}_3^{(b)} \mathcal{A}^{(b)}(01)) \\ & = \alpha_1 \mathcal{S}_3^{(a)} \otimes \mathcal{M}_3^{(b)} (1 - A^{(b)}). \end{aligned}$$

This part should be smaller than the projector onto the subspace, $\mathcal{S}_3^{(a)} \otimes \mathcal{M}_3^{(b)}$. Here, $A^{(b)}$ is defined to be $(T^{(b)}(01) + T^{(b)}(02))/2$ in the same way for the global operator A defined by Eq.(7). The eigenvalues of $A^{(b)}$ in $V_M^{(b)}$ are $1/2$ and $-1/2$. Therefore, we obtain $\alpha_1 \leq 2/3$. Similarly, we obtain $\alpha_4 \leq 2/3$ from the inequality in the subspace $V_M^{(a)} \otimes V_S^{(b)}$.

The inequality $E_1^L + E_2^L \leq 1$ in $V_A^{(a)} \otimes V_M^{(b)}$ takes the form:

$$\begin{aligned} & \alpha_2 (\mathcal{A}_3^{(a)} \otimes \mathcal{M}_3^{(b)} \mathcal{S}^{(b)}(02) + \mathcal{A}_3^{(a)} \otimes \mathcal{M}_3^{(b)} \mathcal{S}^{(b)}(01)) \\ & = \alpha_2 \mathcal{A}_3^{(a)} \otimes \mathcal{M}_3^{(b)} (1 + A^{(b)}) \leq \mathcal{A}_3^{(a)} \otimes \mathcal{M}_3^{(b)}, \end{aligned}$$

which requires that $\alpha_2 \leq 2/3$. In the same way, we obtain $\alpha_3 \leq 2/3$ from the inequality in $V_M^{(a)} \otimes V_A^{(b)}$. Thus all the four coefficients α should be less than or equal to $2/3$.

It is not straightforward to find allowed ranges of β_1 and β_2 from the inequality in $V_M^{(a)} \otimes V_M^{(b)}$. In the space

$V_M^{(a)} \otimes V_M^{(b)}$, we define an operator X to be

$$\begin{aligned} X \equiv & \beta_1 (\mathcal{S}^{(a)}(02)\mathcal{A}^{(b)}(02) + \mathcal{S}^{(a)}(01)\mathcal{A}^{(b)}(01)) \\ & + \beta_2 (\mathcal{A}^{(a)}(02)\mathcal{S}^{(b)}(02) + \mathcal{A}^{(a)}(01)\mathcal{S}^{(b)}(01)), \end{aligned} \quad (47)$$

which is the part of $E_1^L + E_2^L$ contributing to the space $V_M^{(a)} \otimes V_M^{(b)}$. We need to find the greatest eigenvalue of X , since the inequality implies $X \leq \mathcal{M}_3^{(a)} \otimes \mathcal{M}_3^{(b)}$. It should be understood that we are working in subspace $V_M^{(a)} \otimes V_M^{(b)}$, and the projectors $\mathcal{M}_3^{(a)}$ and $\mathcal{M}_3^{(b)}$ will be omitted. In terms of $A^{(p)}$ and $D^{(p)}$, the operator X is expressed as

$$\begin{aligned} X = & \beta (1 - A^{(a)} \otimes A^{(b)} - D^{(a)} \otimes D^{(b)}) \\ & + \delta (1^{(a)} \otimes A^{(b)} - A^{(a)} \otimes 1^{(b)}), \end{aligned}$$

where $\beta = \frac{1}{2}(\beta_1 + \beta_2)$ and $\delta = \frac{1}{2}(\beta_1 - \beta_2)$. In order to diagonalize X , it is convenient to introduce the basis in which $A^{(p)}(p = a, b)$ is diagonal:

$$\begin{aligned} A^{(p)}|m+\rangle &= \frac{1}{2}|m+\rangle, \quad A^{(p)}|m-\rangle = -\frac{1}{2}|m-\rangle, \\ D^{(p)}|m+\rangle &= \frac{\sqrt{3}}{2}|m-\rangle, \quad D^{(p)}|m-\rangle = \frac{\sqrt{3}}{2}|m+\rangle, \end{aligned}$$

where $m = 1, 2, \dots, \dim(V_M^{(p)})/2$. The bipartite state $|m\pm\rangle \otimes |m'\pm\rangle$ for a given set of m and m' will be written as $|\pm\pm\rangle$ for simplicity.

In this basis, two eigenvalues of X are easily found by inspection:

$$\begin{aligned} X(|++\rangle + |--\rangle) &= 0, \\ X(|++\rangle - |--\rangle) &= \frac{3}{2}\beta(|++\rangle - |--\rangle). \end{aligned}$$

States $|+-\rangle$ and $|-+\rangle$ are transformed by X as

$$\begin{aligned} X|+-\rangle &= \left(\frac{5}{4}\beta - \delta\right)|+-\rangle - \frac{3}{4}\beta|-\rangle, \\ X|-+\rangle &= \left(\frac{5}{4}\beta + \delta\right)|-+\rangle - \frac{3}{4}\beta|+-\rangle. \end{aligned}$$

The other two eigenvalues are determined by diagonalizing the 2×2 matrix corresponding to the above transformation and found to be

$$\gamma_{\pm} = \frac{5}{4}\beta \pm \sqrt{\frac{9}{16}\beta^2 + \delta^2}.$$

Of the four eigenvalues, the greatest one is γ_+ . The positivity of E_0^L thus requires that the positive coefficients β_1 and β_2 should satisfy the condition:

$$\frac{5}{4}\beta + \sqrt{\frac{9}{16}\beta^2 + \delta^2} \leq 1, \quad (48)$$

where $\beta = \frac{1}{2}(\beta_1 + \beta_2)$ and $\delta = \frac{1}{2}(\beta_1 - \beta_2)$.

C. Maximum success probability by separable POVM

Now that we have the possible form of the separable POVM E_μ^L and the conditions for the coefficients in it, we can optimize the mean success probability given by

$$p^L = \frac{1}{d_2 d_1} \text{tr}[E_1^L \mathcal{S}(01)]. \quad (49)$$

The trace $\text{tr}[E_1^L \mathcal{S}(01)]$ can be calculated by decomposing the trace into traces in the subsystems as

$$\begin{aligned} & \text{tr}[E_1^L \mathcal{S}(01)] \\ &= \text{tr}\left[E_1^L \left(\mathcal{S}^{(a)}(01)\mathcal{S}^{(b)}(01) + \mathcal{A}^{(a)}(01)\mathcal{A}^{(b)}(01)\right)\right]. \end{aligned}$$

We must calculate many traces in subsystems, for which the following formulas can be used ($p = a, b$):

$$\begin{aligned} & \text{tr}\mathcal{M}_3^{(p)} \mathcal{S}^{(p)}(02)\mathcal{S}^{(p)}(01) = \text{tr}\mathcal{M}_3^{(p)} \mathcal{A}^{(p)}(02)\mathcal{A}^{(p)}(01) \\ &= \frac{1}{2} \text{tr}\left[\mathcal{M}_3^{(p)} (A^{(p)})^2\right] = \frac{1}{8} \dim V_{\mathcal{M}}^{(p)}, \\ & \text{tr}\mathcal{M}_3^{(p)} \mathcal{S}^{(p)}(02)\mathcal{A}^{(p)}(01) = \text{tr}\mathcal{M}_3^{(p)} \mathcal{A}^{(p)}(02)\mathcal{S}^{(p)}(01) \\ &= \frac{1}{2} \text{tr}\left[\mathcal{M}_3^{(p)} (D^{(p)})^2\right] = \frac{3}{8} \dim V_{\mathcal{M}}^{(p)}. \end{aligned}$$

The result is given by

$$\begin{aligned} & \text{tr}[E_1^L \mathcal{S}(01)] \\ &= \frac{3}{8} \left(\alpha_1 \dim V_S^{(a)} \dim V_{\mathcal{M}}^{(b)} + \alpha_2 \dim V_{\mathcal{A}}^{(a)} \dim V_{\mathcal{M}}^{(b)} \right. \\ & \quad \left. + \alpha_3 \dim V_{\mathcal{M}}^{(a)} \dim V_{\mathcal{A}}^{(b)} + \alpha_4 \dim V_{\mathcal{M}}^{(a)} \dim V_S^{(b)} \right) \\ & \quad + \frac{3}{32} (\beta_1 + \beta_2) \dim V_{\mathcal{M}}^{(a)} \dim V_{\mathcal{M}}^{(b)}. \end{aligned}$$

It is clear that we should take the largest possible value $2/3$ for the coefficients α in order to maximize $\text{tr}[E_1^L \mathcal{S}(01)]$. For β_1 and β_2 , note that $\text{tr}[E_1^L \mathcal{S}(01)]$ contains β 's in the form of $\beta_1 + \beta_2$, and we can use the following inequalities:

$$\beta_1 + \beta_2 = 2\beta \leq \frac{5}{4}\beta + \sqrt{\frac{9}{16}\beta^2 + \delta^2} \leq 1. \quad (50)$$

Evidently, $\beta_1 + \beta_2$ takes the maximum value 1 only when $\beta_1 = \beta_2 = 1/2$. The maximum value of $\text{tr}[E_1^L \mathcal{S}(01)]$ with separable POVM is thus given by

$$\begin{aligned} & \text{tr}[E_1^L \mathcal{S}(01)] \\ &= \frac{1}{4} \left(\dim V_S^{(a)} \dim V_{\mathcal{M}}^{(b)} + \dim V_{\mathcal{A}}^{(a)} \dim V_{\mathcal{M}}^{(b)} \right. \\ & \quad \left. + \dim V_{\mathcal{M}}^{(a)} \dim V_{\mathcal{A}}^{(b)} + \dim V_{\mathcal{M}}^{(a)} \dim V_S^{(b)} \right) \\ & \quad + \frac{3}{8} \dim V_{\mathcal{M}}^{(a)} \dim V_{\mathcal{M}}^{(b)}. \end{aligned} \quad (51)$$

On the other hand, we have

$$\text{tr}[E_1 \mathcal{S}(01)] = \frac{1}{4} \dim V_{\mathcal{M}}, \quad (52)$$

for the global POVM element E_1 . Using the relation of dimensions given in Eq.(26), we thus conclude that $\text{tr}[E_1^L \mathcal{S}(01)] < \text{tr}[E_1 \mathcal{S}(01)]$, implying that any separable POVM cannot attain the maximum unambiguous identification probability by the global measurement scheme.

D. LOCC protocol

Thus, the optimal separable POVM element E_1^L is given by Eq.(46) with $\alpha_i = 2/3$ and $\beta_i = 1/2$:

$$\begin{aligned} E_1^L &= \frac{2}{3} \left(\mathcal{S}_3^{(a)} \otimes \mathcal{M}_3^{(b)} \mathcal{A}^{(b)}(02) + \mathcal{A}_3^{(a)} \otimes \mathcal{M}_3^{(b)} \mathcal{S}^{(b)}(02) \right. \\ & \quad \left. + \mathcal{M}_3^{(a)} \mathcal{S}^{(a)}(02) \otimes \mathcal{A}_3^{(b)} + \mathcal{M}_3^{(a)} \mathcal{A}^{(a)}(02) \otimes \mathcal{S}_3^{(b)} \right) \\ & \quad + \frac{1}{2} \left(\mathcal{M}_3^{(a)} \mathcal{S}^{(a)}(02) \otimes \mathcal{M}_3^{(b)} \mathcal{A}^{(b)}(02) \right. \\ & \quad \left. + \mathcal{M}_3^{(a)} \mathcal{A}^{(a)}(02) \otimes \mathcal{M}_3^{(b)} \mathcal{S}^{(b)}(02) \right). \end{aligned} \quad (53)$$

The remaining elements are given as $E_2^L = T(12)E_1^L T(12)$ by the exchange symmetry, and $E_0^L = 1 - E_1^L - E_2^L$ by the completeness of the POVM. We can now show that this separable POVM E_μ^L can be implemented by a LOCC protocol, which is summarized as follows:

- First, Alice and Bob determine the permutation symmetry of their local system: totally symmetric, mixed symmetric, or totally antisymmetric. This is done by a projective measurement with the set of orthogonal projectors $\{\mathcal{S}_3^{(p)}, \mathcal{M}_3^{(p)}, \mathcal{A}_3^{(p)}\}$ of each party $p(= a, b)$. Note that the cases $\mathcal{S}_3^{(a)} \otimes \mathcal{A}_3^{(b)}$ and $\mathcal{A}_3^{(a)} \otimes \mathcal{S}_3^{(b)}$ do not occur, since they would imply that the whole system is totally antisymmetric.
- If their outcome is $\mathcal{S}_3^{(a)} \otimes \mathcal{S}_3^{(b)}$ or $\mathcal{A}_3^{(a)} \otimes \mathcal{A}_3^{(b)}$, Alice and Bob declare an inconclusive result, i.e., 0.
- If one of the two parties $p(= a \text{ or } b)$ finds that his or her local system is totally symmetric, $\mathcal{S}_3^{(p)}$, and the system of the other party $q(\neq p)$ is found to be mixed symmetric, $\mathcal{M}_3^{(q)}$, then party q performs a POVM measurement:

$$\begin{aligned} e_1 &\equiv \frac{2}{3} \mathcal{M}_3^{(q)} \mathcal{A}^{(q)}(02), \\ e_2 &\equiv \frac{2}{3} \mathcal{M}_3^{(q)} \mathcal{A}^{(q)}(01), \\ e_0 &\equiv \frac{1}{3} \mathcal{M}_3^{(q)} \left(1 + 2A^{(q)} \right). \end{aligned}$$

Note that the set $\{e_1, e_2, e_0\}$ is a POVM since $e_\mu \geq 0$, ($\mu = 0, 1, 2$) and $\sum_\mu e_\mu = \mathcal{M}_3^{(q)}$. The final identification result by Alice and Bob is the measurement outcome $\mu(= 0, 1, 2)$ of party q .

- If one of the two parties $p(= a \text{ or } b)$ finds that his or her local system is totally antisymmetric, $\mathcal{A}_3^{(p)}$,

and the system of the other party $q (\neq p)$ is found to be mixed symmetric, $\mathcal{M}_3^{(q)}$, then party q performs a POVM measurement:

$$\begin{aligned} e'_1 &\equiv \frac{2}{3} \mathcal{M}_3^{(q)} \mathcal{S}^{(q)}(02), \\ e'_2 &\equiv \frac{2}{3} \mathcal{M}_3^{(q)} \mathcal{S}^{(q)}(01), \\ e'_0 &\equiv \frac{1}{3} \mathcal{M}_3^{(q)} (1 - 2A^{(q)}). \end{aligned}$$

It is easily verified that the set $\{e'_1, e'_2, e'_0\}$ is also a POVM in $V_{\mathcal{M}}^{(q)}$, and the final identification result of Alice and Bob is chosen to be the measurement outcome $\mu (= 0, 1, 2)$ of party q .

- Finally, when the total system is found to be in $V_{\mathcal{M}}^{(a)} \otimes V_{\mathcal{M}}^{(b)}$, one of the two parties, say Alice, performs the following POVM measurement:

$$\begin{aligned} e_{11} &\equiv \frac{1}{2} \mathcal{M}_3^{(a)} \mathcal{A}^{(a)}(02), \quad e_{12} \equiv \frac{1}{2} \mathcal{M}_3^{(a)} \mathcal{S}^{(a)}(02), \\ e_{21} &\equiv \frac{1}{2} \mathcal{M}_3^{(a)} \mathcal{A}^{(a)}(01), \quad e_{22} \equiv \frac{1}{2} \mathcal{M}_3^{(a)} \mathcal{S}^{(a)}(01). \end{aligned}$$

It is evident that the above set $\{e_{a_1 a_2}\}_{a_1, a_2=1,2}$ forms a POVM in $V_{\mathcal{M}}^{(a)}$. If Alice's outcome a_1 is equal to 1, Bob performs a projective measurement by the set of orthogonal projectors $\{f_b\}_{b=1,2}$:

$$f_1 \equiv \mathcal{M}_3^{(b)} \mathcal{S}^{(b)}(02), \quad f_2 \equiv \mathcal{M}_3^{(b)} \mathcal{A}^{(b)}(02),$$

otherwise, by the set of orthogonal projectors $\{f'_b\}_{b=1,2}$:

$$f'_1 \equiv \mathcal{M}_3^{(b)} \mathcal{S}^{(b)}(01), \quad f'_2 \equiv \mathcal{M}_3^{(b)} \mathcal{A}^{(b)}(01).$$

The final identification result μ of Alice and Bob is Alice's result a_1 , if Alice's outcome a_2 coincides with Bob's outcome b . Otherwise, the final result is the inconclusive one, i.e., 0.

Substituting explicit dimensions in Eq.(51), we obtain the optimal success probability with the LOCC protocol:

$$p_{\max}^L = \frac{1}{36d_a d_b (d_a d_b + 1)} (11d_a^2 d_b^2 + d_a^2 + d_b^2 - 13), \quad (54)$$

whereas the globally attainable success probability of Eq.(41) in terms of dimensions d_a and d_b is given by

$$p_{\max} = \frac{1}{3d_a d_b} (d_a d_b - 1). \quad (55)$$

Although there is a finite gap between p_{\max} and p_{\max}^L as shown before, the numerical difference is not very large. For example, in the case of a two-qubit bipartite system ($d_a = d_b = 2$), the optimal LOCC protocol gives $p_{\max}^L = 19/80$, whereas the globally attainable probability is given by $p_{\max} = 1/4$. The difference is only $1/80$. In the limit of d_a and d_b going to infinity, we find that p_{\max}^L approaches $11/36$ and p_{\max} approaches $1/3$ with the difference $1/36$.

VI. CONCLUDING REMARKS

It is known that two bipartite pure states can be optimally discriminated within a LOCC scheme if classical knowledge of the two states are available. In this paper, we investigated the identification problem of two bipartite pure states, where no classical knowledge of the reference states is given but only a copy of each reference state is available. The two reference states are independently and randomly chosen from the state space in a unitary invariant way.

In the case of minimum-error identification, we found that the optimal identification can be done locally. This is true for any prior probabilities, but we assumed that the number N of copies of each state is 1. In the limit of large N , the identification problem reduces to the standard discrimination problem. This is because one can obtain complete classical information on the reference states by performing a tomographical measurement on infinitely many copies of them. We note that an infinite number of copies of a bipartite state is not equivalent to perfect classical knowledge of the state. In general, possession of an unlimited number of copies is more advantageous than mere classical knowledge of the state, because two parties sharing the states can utilize a quantum channel. However, for the identification problems with $N = \infty$ considered in this paper, it turns out not to be an advantage. This is because the discrimination of two known pure states can be performed optimally by LOCC and no scheme can perform better than the optimal global protocol. Therefore, we can say that the pure-state identification can be optimally performed by means of LOCC when $N = 1$ and $N = \infty$. We conjecture that this is also true for arbitrary N .

On the other hand, we have demonstrated that any LOCC scheme for the unambiguous identification of two bipartite pure states cannot attain the maximum success probability achieved by the global measurement at least for the case of equal prior probabilities. This contrasts remarkably with the results for the standard discrimination and the minimum-error identification. Our result provides an example of nonlocality in distinguishing two pure states. It is an interesting problem to study the unambiguous identification of bipartite pure states when the number of copies of the reference states is finite but greater than 1.

Just for completeness, we comment on the $N = 0$ case, where no measurement (global or LOCC) on the input state can improve the success probability. The maximum success probabilities are $\max\{\eta_1, \eta_2\}$ and 0 for the minimum-error and the unambiguous identification problems, respectively.

In this paper we assumed that the two reference states are independently and uniformly distributed in the state space. This is a sensible assumption to make for two completely unknown states. However, we note that the optimal identification protocol depends on the distribution of the reference states and whether the LOCC scheme

attains global optimality may also change for different distributions.

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